

Asymptotics, frequency modulation, and low regularity ill-posedness for canonical defocusing equations

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after

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We consider the cubic one-dimensional defocusing nonlinear Schrödinger equation (NLS):

$$\begin{cases} -iu_t + \Delta u = |u|^2 u \\ u(0, x) = u_0(x) \end{cases} \quad u_0 \in H_x^s(\mathbb{R}). \quad (\text{NLS})$$

We consider solutions in the space $C^0([0, T], H^s(\mathbb{R}))$ that are fixed points of the Duhamel map:

$$u(t) = U(t)u_0 + i \int_0^t U(t-s)|u|^2 u(s) \, ds,$$

with $U(t) = e^{-it\Delta}$ defined as $U(t)u_0 := \mathcal{F}^{-1}(e^{it\xi^2}\widehat{u_0})$, where \mathcal{F} is the Fourier Transform.

1 Well-posedness

We will say that our Initial Value Problem is

locally well-posed in H^s if for every ball B_R of radius $R > 0$ exists a time $T = T(R) > 0$ such that the solution operator

$$S: B_R = \{u_0 \in H^s : \|u_0\|_{H^s} < R\} \rightarrow C^0([0, T], H^s(\mathbb{R}))$$

is uniformly continuous.

globally well-posed if T can be arbitrarily large independent of R .

1.1 Previous result

Theorem 1 (Tsutsumi, 1987). *The defocusing (NLS) is globally well-posed in $H^s(\mathbb{R})$ for $s \geq 0$.*

The main result of this talk is:

Theorem 2. *The cubic one-dimensional defocusing (NLS) is not locally well-posed in $H^s(\mathbb{R})$ for $s < 0$.*

2 Symmetries

Let $u(t, x)$ be a solution of (NLS), then are solutions as well:

Symmetry		Invariant norm
Scaling	$u^\lambda(t, x) = \frac{1}{\lambda} u\left(\frac{t}{\lambda^2}, \frac{x}{\lambda}\right)$	$\dot{H}_x^{-1/2}$
Galilean	$G_v(u)(t, x) = e^{itv^2} e^{ixv} u(t, x + 2vt)$	L_x^2

$$H^s(\mathbb{R}) := \left\{ u \in \mathcal{S}'(\mathbb{R}) \mid (1 + |\xi|^2)^{s/2} \widehat{u}(\xi) \in L^2(\mathbb{R}) \right\}$$

2.1 Free Schrödinger and pseudo-conformal transformation

Consider the Free Schrödinger Equation for forward time

$$\begin{cases} -iu_t + \Delta u = 0 & t > 0 \\ u(0, x) = u_0(x). \end{cases} \quad (\text{FSE})$$

Let $u(t, x)$ be a solution of (FSE). We introduce the *pseudo-conformal* transformation by setting:

$$v(s, y) := \text{pc}(u) = s^{-1/2} \exp(iy^2/4s) u\left(\frac{1}{s} - 1, \frac{y}{s}\right).$$

This transformation is associated to the following change of variables:

$$(y, s) = \left(\frac{x}{1+t}, \frac{1}{1+t} \right), \quad (x, t) = \left(\frac{y}{s}, \frac{1}{s} - 1 \right).$$

The transformed $v(s, y) := \text{pc}(u)$ solves the backwards Free Schrödinger for time :

$$\begin{cases} iv_s + \Delta v = 0 & s \in (0, 1] \\ v(1, y) = v_1(y). \end{cases} \quad (\text{bFSE})$$

We now can solve this problem for any time $s < 1$, and in particular the solution extends continuously to $s = 0$:

$$v(s, y) \xrightarrow[L^2]{s \rightarrow 0^+} v(0, y) =: \varphi(y).$$

Using the inverse transformation $u = \text{pc}^{-1}(v)$ we see that:

$$u(t, x) \stackrel{t \rightarrow +\infty}{\approx} (1+t)^{-1/2} \exp(-ix^2/4(1+t)) \varphi\left(\frac{x}{1+t}\right).$$

3 Backwards Nonlinear Schrödinger equation

Applying the pc transformation to solutions of (NLS) we obtain a new initial value problem for $v(s, y) := \text{pc}(u)$:

$$\begin{cases} iv_s + \Delta v = s^{-1}|v|^2 v & s \in (0, 1] \\ v(1, y) = v_1(y) \end{cases} \quad (\text{bNLS})$$

Dropping $\partial_{yy}v$, we can explicitly solve the ODE:

$$\begin{cases} iv_s^{[w]} = s^{-1}|v^{[w]}|^2v^{[w]} \\ v(1, y) = w(y) \end{cases} \quad (\text{ODE})$$

finding the solution

$$v^{[w]}(s, y) = w(y) \exp(-i|w(y)|^2 \log(s)).$$

We notice that as $s \rightarrow 0^+$:

$$|v_s^{[w]}| \sim s^{-1}, \quad |\Delta v^{[w]}| \sim (\log s)^2, \quad \left|s^{-1}|v^{[w]}|^2v^{[w]}\right| \sim s^{-1},$$

and since $|\log s| \ll s^{-1}$, neglecting the term $\Delta v^{[w]}$ is reasonable and $v^{[w]}$ turns out to be “close” to a solution of (bNLS).

To make this precise we introduce a weighted Sobolev norm:

$$\|v\|_{\mathbf{H}_y^{k,k}} := \sum_{\substack{\alpha, \beta \geq 0 \\ \alpha + \beta \leq k}} \left\| y^\alpha \partial_y^\beta v \right\|_{L_y^2}.$$

Consider the ball $B_\epsilon = \{\|u\|_{\mathbf{H}^{k+2,k+2}} < \epsilon\}$, where $\epsilon \ll 1$. Let be $k \in \mathbb{N}, k \geq 5$.

Lemma 1. *For all $w \in B_\epsilon$ exists $v_1 \in \mathbf{H}^{k,k}$ such that the unique solution $v^{\langle w \rangle}$ to (bNLS) in $\mathbf{H}^{k,k}$ with initial data v_1 satisfies*

$$\left\| v^{\langle w \rangle}(s) - v^{[w]}(s) \right\|_{\mathbf{H}^{k,k}} \lesssim \epsilon s (1 + |\log s|)^C \quad \text{for all } s \in (0, 1]$$

Furthermore, the map from B_ϵ to $C^0((0, 1], \mathbf{H}^{k,k})$, $w \mapsto v^{\langle w \rangle}$ is Lipschitz up to the end time $s = 0$.

3.1 Decoherence in L^2

We fix w (e.g. $w(y) = \epsilon e^{-y^2}$, so $w \approx \epsilon$). For $a \in \mathbb{R}$ we can take the initial value $aw(y)$. We have this:

Lemma 2 (decoherence property). *If $a, a' \lesssim 1$ and $a \neq a'$, then*

$$\limsup_{s \rightarrow 0^+} \left\| v^{[aw]}(s) - v^{[a'w]}(s) \right\|_{L_y^2} \gtrsim (|a| + |a'|) \|w\|_{L_y^2}.$$

The last two Lemmas together imply

$$\limsup_{s \rightarrow 0^+} \left\| v^{\langle aw \rangle}(s) - v^{\langle a'w \rangle}(s) \right\|_{L_y^2} \gtrsim \epsilon.$$

Using the inverse pc we may construct a family of solutions $u^{\langle aw \rangle} = \text{pc}^{-1}(v^{\langle aw \rangle})$, and $u^{[aw]} = \text{pc}^{-1}(v^{[aw]})$ such that:

- $\|u^{\langle aw \rangle}(0)\|_{H^k} \lesssim \|v^{\langle aw \rangle}(1)\|_{\mathbf{H}^{k,k}} \lesssim \epsilon$;
- $\left\| u^{\langle aw \rangle}(0) - u^{\langle a'w \rangle}(0) \right\|_{H^k} \lesssim \left\| v^{\langle aw \rangle}(1) - v^{\langle a'w \rangle}(1) \right\|_{\mathbf{H}^{k,k}} \lesssim \epsilon |a - a'|$;
- $\left\| u^{\langle aw \rangle}(t) - u^{[a'w]}(t) \right\|_{H^k} \lesssim \epsilon (1+t)^{-1} (1 + |\log(1+t)|)^C$;
- $\limsup_{t \rightarrow +\infty} \left\| u^{\langle aw \rangle}(t) - u^{\langle a'w \rangle}(t) \right\|_{L^2} = \limsup_{s \rightarrow 0^+} \left\| v^{\langle aw \rangle}(s) - v^{\langle a'w \rangle}(s) \right\|_{L^2} \gtrsim \epsilon$.

4 Ill-posedness

To disprove uniform continuity of $S: B_R \subset H^s(\mathbb{R}) \rightarrow C^0([0, T], H^s(\mathbb{R}))$, for $s \in (-\frac{1}{2}, 0)$, we will show that for any $T > 0$ there exists $\epsilon > 0$ such that for any $\delta > 0$ there are two solutions ϕ, ϕ' such that:

- $\phi(0), \phi'(0) \in B_R$,
- $\|\phi(0) - \phi'(0)\|_{H^s} \lesssim \delta$,
- $\sup_{[0, T]} \|\phi(t) - \phi'(t)\|_{H^s} \gtrsim \epsilon$.

We take $N \gg 1$ and consider:

$$\phi^{(a)}(t, x) := G_N(u^\lambda)(t, x) = \frac{1}{\lambda} e^{itN^2} e^{ixN} u^{(aw)}\left(\frac{t}{\lambda^2}, \frac{x + 2Nt}{\lambda}\right).$$

We need a lemma to control the H^s norm of $\phi^{(a)}$ by the H^k norm from above and by the L^2 norm from below:

Lemma 3. *Let be $k \in \mathbb{N}$, $u \in H^k$, $N \geq 1$, $\lambda \in \mathbb{R}^+$, and let be $\phi^{(a)}$ defined as above. Then*

i) *For $s < 0$, $k \geq |s|$, whenever $1 \leq \lambda N^{1+(s/k)}$, we have:*

$$\|\phi\|_{H^s} \lesssim \lambda^{-\frac{1}{2}} N^s \|u\|_{H^k},$$

ii) *For every $u \in H^k$ exists a constant $C_u < \infty$ such that when $\lambda N \geq C_u$ we have:*

$$\|\phi\|_{H^s} \gtrsim \lambda^{-\frac{1}{2}} N^s \|u\|_{L^2}$$

Taking $a, a' \lesssim 1$, consider $\phi^{(a)}$ and $\phi^{(a')}$. Applying the Lemma and choosing $\lambda = N^{2s}$ these two solutions disprove the uniform continuity of S , in fact:

- a) $\|\phi^{(a)}(0)\|_{H^s} \lesssim \|u^{(aw)}(0)\|_{H^k} \lesssim \epsilon$;
- b) $\|\phi^{(a)}(0) - \phi^{(a')}(0)\|_{H^s} \lesssim \epsilon |a - a'|$;
- d) $\limsup_{t \rightarrow +\infty} \|\phi^{(a)}(t) - \phi^{(a')}(t)\|_{H^s} = \limsup_{t \rightarrow +\infty} \|u^{(aw)}(t) - u^{(a'w)}(t)\|_{L^2} \gtrsim \epsilon$.

Thus there exists t_0 such that $\|u^{(aw)}(t_0) - u^{(a'w)}(t_0)\|_{L^2} \gtrsim \epsilon$. Let $t = \lambda^2 t_0$. Making N arbitrarily large, $\lambda = N^{2s}$ becomes arbitrarily small, so $t \in [0, T)$ for any time $T > 0$.